

A TWO-FREQUENCY METHOD FOR IMPEDANCE MEASUREMENT IN MICROWAVE
FREQUENCY-HALVING NETWORKS

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ABSTRACT

The method described removes a major difficulty in the design of wideband frequency-halvers: the accurate characterization of the "pumped" input and output impedances at the plane of the varactors. The procedure permits the use of realistic input and output power levels with the input at twice the output frequency.

INTRODUCTION

The optimum design of wideband frequency halvers requires that the "pumped" input and output varactor impedances Z_{in} and Z_{out} be known *a priori*. Since the halver is a threshold device [1,2] these impedances exist only for appropriate values of input power P_{in} and input frequency 2ω . Therefore $Z_{in}(2\omega)$ and $Z_{out}(\omega)$ can only be measured in a functioning halver circuit. This dilemma can be resolved thus:

- (a) Assuming the validity of the "pumped" varactor equations of Penfield and Rafuse [3], determine estimates $\tilde{Z}_{in}(2\omega)$ and $\tilde{Z}_{out}(\omega)$.
- (b) Use \tilde{Z}_{in} and \tilde{Z}_{out} to design a prototype halver, e.g. with the aid of COMPACT [4].
- (c) Measure the external impedances Z_{in}^{ext} and Z_{out}^{ext} of the resulting two-port.
- (d) Knowing the halver network, "de-embed" the varactor impedances Z_{in} and Z_{out} .
- (e) Compare Z_{in} and Z_{out} with \tilde{Z}_{in} and \tilde{Z}_{out} .
- (f) Iterate as necessary.

Impedance measurement by conventional load-pull techniques requires determination of large-signal S-parameters. At a specific input frequency 2ω and power P_{in} , the halver can be characterized by

$$\begin{bmatrix} b_1(2\omega) \\ b_2(\omega) \end{bmatrix} = \begin{bmatrix} S_{11} & S_{12} \\ S_{21} & S_{22} \end{bmatrix} \cdot \begin{bmatrix} a_1(2\omega) \\ a_2(\omega) \end{bmatrix}, \quad (1)$$

see Fig. 1. For a balanced halver, the port signals have minimum harmonic content [2]. From (1):

$$\begin{aligned} S_{11} &= \frac{b_1(2\omega)}{a_1(2\omega)} \Bigg|_{a_2(\omega)=0}, & S_{12} &= \frac{b_1(2\omega)}{a_2(\omega)} \Bigg|_{a_1(2\omega)=0} \\ S_{21} &= \frac{b_2(\omega)}{a_1(2\omega)} \Bigg|_{a_2(\omega)=0}, & S_{22} &= \frac{b_2(\omega)}{a_2(\omega)} \Bigg|_{a_1(2\omega)=0}. \end{aligned} \quad (2)$$

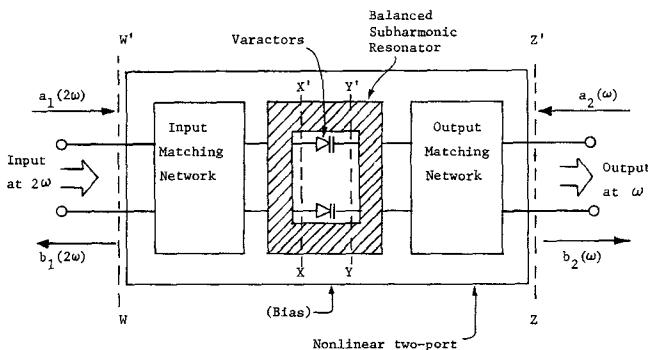


Fig. 1 Definition of 2-frequency S-matrix and reference planes.

A network analyser (NA) can be used to measure S_{11} at 2ω with the output terminated, but not the trans-frequency parameters S_{12} and S_{21} (which are of minor concern here). Determination of S_{22} according to (2) is impossible because a reverse signal $a_2(\omega)$ has to be applied to the output port and the resulting reflected signal $b_2(\omega)$ measured, the input being terminated ($a_1=0$). However, this removes the drive power, resulting in cessation of halving.

TWO-FREQUENCY MEASUREMENT TECHNIQUE

A solution employs a two-frequency variant of Mazumder's synthetic load technique [5]. From (1):

$$\frac{b_1(2\omega)}{a_1(2\omega)} = S_{11} + S_{12} \frac{a_2(\omega)}{a_1(2\omega)}, \quad (3)$$

$$\frac{b_2(\omega)}{a_2(\omega)} = S_{21} \frac{a_1(2\omega)}{a_2(\omega)} + S_{22}. \quad (4)$$

For fixed $|a_1|$ and $|a_2|$, variation of $\angle a_1 - \angle a_2$ generates circles in the complex planes of b_1/a_1 and b_2/a_2 , see Fig.2. The centres of these circles correspond to S_{11} and S_{22} respectively. NA measurements permit b_1/a_1 and b_2/a_2 to be found by varying $\angle a_1 - \angle a_2$ with a phase-shifter.

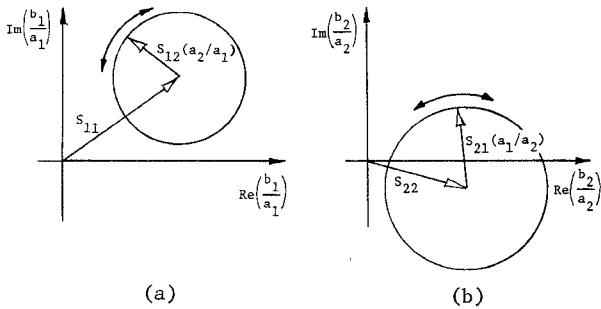


Fig. 2 Circular loci in: (a) b_1/a_1 -plane at 2ω , (b) b_2/a_2 -plane at ω .

Measurement Procedure

For each input frequency, three steps are taken:

(i) Reflection coefficients Γ_{mj} ($j = 1, 2, 3$) are measured at reference planes WW' and ZZ' (see Fig.1) as described below.

Fig.3 shows a set-up which applies $a_1(2\omega)$ to the halver input and $a_2(\omega)$ to its output. Attenuator A_1 sets $|a_1|$, while A_2 and P set $|a_2|$ and $\angle a_1 - \angle a_2$ respectively. For accuracy $|a_2|^2$ should be such that $|b_2|^2 - |a_2|^2$ approximates the normal P_{out} to a passive load.

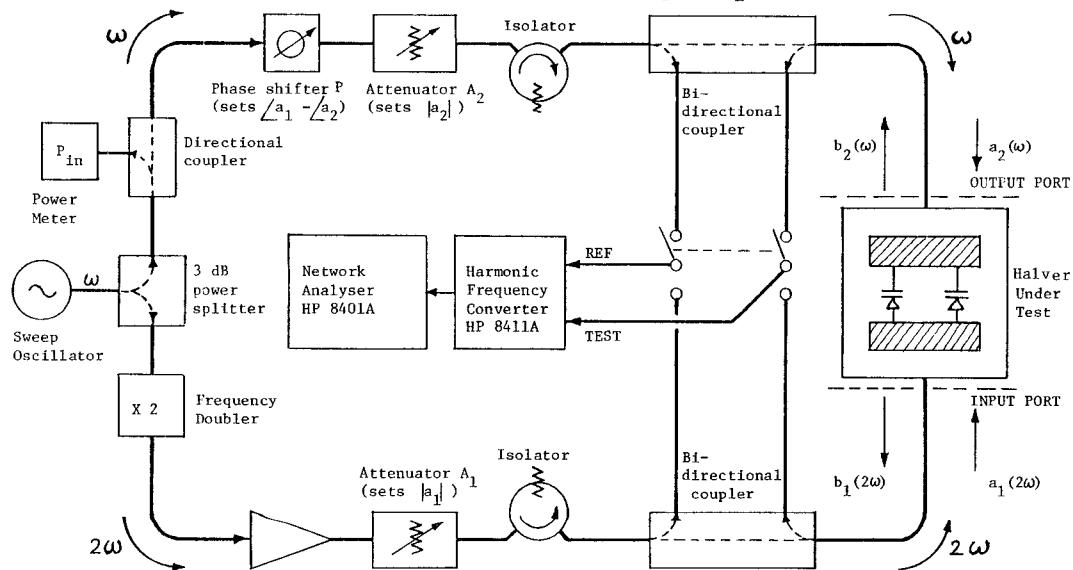


Fig. 3 Two-frequency test set-up.

| | Measured value | True value |
|--------------------|----------------|--------------------|
| Matched load Z_o | Γ_{m1} | $\Gamma_{t1} = 0$ |
| Open circuit | Γ_{m2} | $\Gamma_{t2} = 1$ |
| Short circuit | Γ_{m3} | $\Gamma_{t3} = -1$ |

Hence complex calibration coefficients K, L, M are determined satisfying the bilinear transformation [6]

$$\Gamma_{mj} = (K\Gamma_{tj} + L)/(M\Gamma_{tj} + 1) . \quad (5)$$

Thus for the measured Γ_{mi} ($i = 1, 2$), here representing b_i/a_i , the true values are given by

$$\Gamma_{ti} = (L - \Gamma_{mi})/(M\Gamma_{mi} - K) . \quad (6)$$

(ii) For both b_1/a_1 and b_2/a_2 measurements, Γ_{mk} ($k = 1, 2, 3$) are found for three settings of $\angle a_1 - \angle a_2$. Extraction of the true values Γ_{tk} , according to (6), locates three points on one of the circles of Fig.2, the centres of which correspond to S_{11} and S_{22} . The external halver impedances are then

$$Z_{in}^{ext} = Z_o(1+S_{11})/(1-S_{11}), Z_{out}^{ext} = Z_o(1+S_{22})/(1-S_{22}) . \quad (7)$$

(iii) The varactor embedding being known, reflection coefficients Γ_{in} and Γ_{out} at planes XX' and YY' (Fig.1) can be found using a COMPACT[4] de-embedding file. The large-signal "pumped" varactor impedances Z_{in} , Z_{out} are then obtained from Γ_{in} and Γ_{out} .

RESULTS

The external impedances Z_{in}^{ext} and Z_{out}^{ext} of an experimental halver were measured for $P_{in} = 16.4\text{dBm}$ and $|b_2|^2 - |a_2|^2 = -3\text{ dBm}$, the input frequency being

varied over the 3.5 - 7.0 GHz octave. Fig. 5 compares the de-embedded interior impedances Z_{in} and Z_{out} with theoretical estimates given by the Penfield and Rafuse equations[3] modified to account for parasitics. For two varactors in parallel at 2ω :

$$\tilde{Z}_{in}(2\omega) = \frac{1}{2} \left[r_s + \frac{m_1^2}{2m_2} \cdot \frac{(1+V_V/\phi)^{\gamma}}{2\omega C_j(0)} \right] - j \frac{1}{2} \left[\frac{1}{2\omega C_j(0)} - 2\omega L_s \right],$$

and for two varactors in series at ω :

$$\tilde{Z}_{out}(\omega) = 2 \left[m_2 \frac{(1+V_B/\phi)^{\gamma}}{\omega C_j(0)} - r_s \right] - j 2 \left[\frac{1}{\omega C_j(0)} - \omega L_s \right]. \quad (9)$$

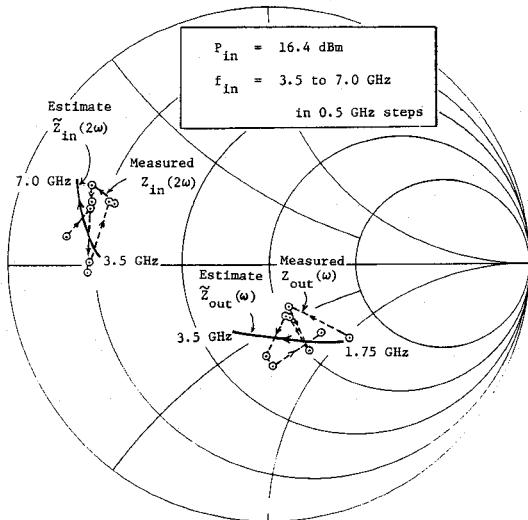


Fig. 5 Comparison between measured $Z_{in}(2\omega)$ and $Z_{out}(\omega)$ of the "pumped" varactors and the estimates $\tilde{Z}_{in}(2\omega)$ and $\tilde{Z}_{out}(\omega)$. Varactors: Frequency Sources, GHZ Division, type GC-1504.

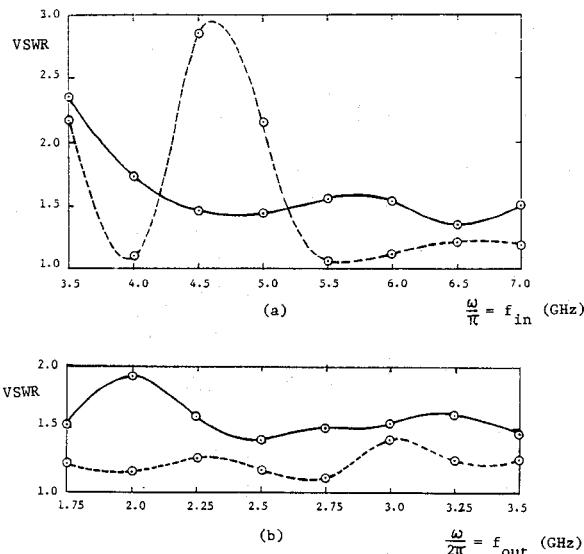


Fig. 6. Measured (-) and predicted (--) VSWRs (a) at the input plane WW' ; (b) at the output plane ZZ' .

Here $C_j(0)$ is the junction capacitance at zero bias, ϕ is the contact potential and γ is the capacitance law exponent. The varactor series resistance and inductance are r_s and L_s respectively. The modulation ratios at the input and output frequencies are taken as $m_1 \approx 0.2$, $m_2 \approx 0.1$, corresponding to maximum P_{out} conditions [3].

Fig. 5 shows satisfactory agreement between the measured varactor impedances $Z_{in}(2\omega)$, $Z_{out}(\omega)$ and the estimates $\tilde{Z}_{in}(2\omega)$, $\tilde{Z}_{out}(\omega)$. In this case a single iteration suffices.

Fig. 6 compares the measured and predicted external VSWRs of the complete halver. The discrepancy in Fig. 6(a) is caused by input circuit tuning adjustments not included in the calculation. The measured and calculated output VSWRs, however, are in good agreement, see Fig. 6(b), since the input circuit was not altered in any way.

CONCLUSIONS

It has been demonstrated that a "two-frequency" variant of Mazumder's method provides an effective and accurate approach to the determination of the "pumped" input and output impedances of a frequency halver, which are required for design centering. A similar approach could be used for doubler design; in that case a halver would replace the doubler in the test set-up.

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